Expansion of the CVBEM into a series using intelligent fractals (IFs)<br>T.V. Hromadka II ${ }^{a}$, R.J. Whitley ${ }^{b}$<br>${ }^{a}$ Department of Mathematics, California State University, Fullerton, CA 92634, USA ${ }^{b}$ Department of Mathematics, University of California, Irvine, CA 92717, USA

## ABSTRACT

In this paper, the development of triangular fractals that geometrically sum into an area whose boundary is a function, of a specific type, is used to expand the Complex Variable Boundary Element Method (or CVBEM) into a series. The entire approximation effort can be written as a sum of Cauchy integrals of incremental changes in basis functions.

## BACKGROUND

The Complex Variable Boundary Element Method, or CVBEM, has been the subject of several papers and books (e.g., Hromadka and Lai, 1987; Hromadka, 1993). The basis of the CVBEM is the use of the Cauchy integal equation to develop approximations of two-dimensional boundary
value problems involving the Laplace and Poisson equations. A property of the CVBEM is that the resulting approximation function, $\widehat{\omega}(z)$, is analytic in the simply connected domain, $\Omega$, and continuous on the problem boundary, $\Gamma$. Thus, $\hat{\omega}(z)=\hat{\phi}(z)+\hat{i}(z)$, where $\hat{\phi}(z)$ and $\hat{\psi}(z)$ are the potential and stream functions, respectively, and both functions satisfy the Laplace equation in $\Omega$.

Let $\omega(z)=\phi(x, y)+i \psi(x, y)$ be analytic on $\Gamma \cup \Omega$, where $\Omega$ is a simply connected domain enclosed by the simple closed boundary $\Gamma$. Since $\omega$ is analytic, $\phi$ and $\psi$ are related by the Cauchy-Reimann equations

$$
\begin{equation*}
\frac{\partial \phi}{\partial x}=\frac{\partial \psi}{\partial y} \text { and } \frac{\partial \phi}{\partial y}=-\frac{\partial \psi}{\partial x} \tag{1}
\end{equation*}
$$

and thus both functions satisfy the two-dimensional Laplace equation in $\Omega$, namely

$$
\begin{equation*}
\frac{\partial^{2} \phi}{\partial x^{2}}+\frac{\partial^{2} \phi}{\partial y^{2}}=0 \text { and } \frac{\partial^{2} \psi}{\partial x^{2}}+\frac{\partial^{2} \psi}{\partial y^{2}}=0 \tag{2}
\end{equation*}
$$

The Cauchy integral theorem states that

$$
\begin{equation*}
\omega(z)=\frac{1}{2 \pi i} \int_{\Gamma} \frac{\omega(\zeta) \mathrm{d} \zeta}{\zeta-z}, z \in \Omega, z \notin \Gamma \tag{3}
\end{equation*}
$$

Let the boundary $\Gamma$ be a polygonal line composed of V vertices. Define nodal points with complex coordinates $z_{j}, j=1, \ldots, m$ on $\Gamma$ such that $m>V$. Nodal points include all the vertices of $\Gamma$ and are numbered in a counter-clockwise direction. Let complex boundary element $\Gamma_{\mathrm{j}}$ be the straight m
line segment joining $z_{j}$ and $z_{j+1}$ so that $\quad \Gamma=\bigcup \Gamma_{j}$. The j=1
CVBEM defines a continuous global trial function, $G(z)$, by

$$
\begin{equation*}
\left.\mathrm{G}(\mathrm{z})=\sum_{\mathrm{j}=1}^{\mathrm{m}} \mathrm{~N}_{\mathrm{j}}(\mathrm{z}){\overline{\left(\phi_{\mathrm{j}}\right.}}_{\mathrm{j}}+\mathrm{i} \bar{\psi}_{\mathrm{j}}\right) \tag{4}
\end{equation*}
$$

where,

$$
N_{j}(z)=\left\{\begin{array}{cl}
\frac{z-z_{j-1}}{z_{j}-z_{j-1}} & z \in \Gamma_{j-1}  \tag{5}\\
0 & z \notin \Gamma_{j} \cup \Gamma_{j-1} \\
\frac{z_{j+1}-z}{z_{j+1}-z_{j}} & z \in \Gamma_{j}
\end{array}\right.
$$

and where $\bar{\phi}_{\mathrm{j}}$ and $\bar{\psi}_{\mathrm{j}}$ are nodal values of the two conjugate components, evaluated at $z_{j}$. An analytic approximation is then determined by

$$
\begin{equation*}
\hat{\omega}(\mathrm{z})=\frac{1}{2 \pi \mathrm{i}} \int_{\Gamma} \frac{\mathrm{G}(\zeta) \mathrm{d} \zeta}{\zeta-\mathrm{z}}, z \in \Omega, z \notin \Gamma \tag{6}
\end{equation*}
$$

If we let $q_{j, j+1}$ be the difference in the polar coordinate angles defined by nodal point coordinates $z_{j+1}$ and $z_{0}$, and $z_{j}$ and $z_{0}$, for $z_{0}$ in $\Omega$, then

$$
\begin{equation*}
\widehat{\omega}\left(z_{0}\right)=\sum_{j=1}^{m}\left[\omega_{j+1}\left(z_{0}-z_{j}\right)-\omega_{j}\left(z_{o}-z_{j+1}\right)\right] \frac{H_{j}}{z_{j+1}-z_{j}} \tag{7}
\end{equation*}
$$

where $\omega_{j}$ and $\omega_{j+1}$ are nodal values at coordinates $z_{j}$ and $z_{j+1}$, and

$$
\begin{equation*}
H_{j}=\ln \left|\frac{z_{j+1}-z_{o}}{z_{j}-z_{o}}\right|+i q_{j, j+1} \tag{8}
\end{equation*}
$$

Since usually only one of the two specified nodal values $\left(\bar{\phi}_{j}, \bar{\psi}_{j}\right)$ is known at each $z_{j}, j=1, \ldots, m$, values for the unknown nodal values must be estimated as part of the CVBEM approach. It is noted that certain continuity requirements are necessary in the complex logarithm terms.

## FRACTAL TRIAL FUNCTION BASIS

The continuous global trial function, $G(\zeta)$, is replaced by a rewriting that is analogous to the triangle fractals used in graphical displays.

In our case, we use the symbol $\Delta_{\mathrm{i}}(\mathrm{z})$ in describing the incremental change between the value of a straightline interpolation between consecutive nodal point values $\phi_{i}$ and $\phi_{j}$, at $x=x_{k}$ and the true value of $\phi$ at $x=x_{k}$, denoted by $\phi_{\mathrm{k}}$.

Given coordinates $z_{i}, z_{j}, z_{k}$ for nodes $i, j, k$, respectively,

$$
\Delta_{\mathrm{i}}^{\mathrm{k}}(\mathrm{z}) \equiv\left\{\begin{array}{cl}
0 ; & \begin{array}{l}
z \notin \text { the boundary element } \\
\text { containing nodes } i, j .
\end{array}  \tag{9}\\
\frac{z-z_{i}}{z_{k}-z_{i}} ; & z \text { between nodes } i, k \text { and } z \in \Gamma . \\
\frac{z_{j}-z}{z_{j}-z_{k}} ; & z \text { between nodes } k, j \text { and } z \in \Gamma .
\end{array}\right.
$$

Hereafter, $\Delta_{\mathrm{i}}^{\mathrm{j}}$ (z) will be simply written as $\Delta_{\mathrm{i}}^{\mathrm{j}}$ as it is understood that a function of $z$ is involved. Straightline interpolation between nodes $i, j$ gives an estimate, $\phi_{k}$, at $z_{k}$ of (see Fig. 1)

$$
\begin{equation*}
\hat{\phi}\left(z_{\mathrm{k}}\right)=\phi_{\mathrm{i}}\left(\frac{z_{\mathrm{j}}-z_{\mathrm{k}}}{z_{\mathrm{j}}-z_{\mathrm{i}}}\right)+\phi_{\mathrm{j}}\left(\frac{z_{\mathrm{k}}-z_{\mathrm{i}}}{z_{\mathrm{j}}-z_{\mathrm{i}}}\right) \tag{10}
\end{equation*}
$$

which would be the value of the global trial function at $\mathrm{z}_{\mathrm{k}}$, $G\left(z_{k}\right)$, for the case where node $k$ is already a node used in $G(\zeta)$. Adding the node $z_{k}$ to $G(\zeta)$ simply adds the incremental contribution of $\phi_{k}$,

$$
\begin{equation*}
\mathrm{G}(\zeta) \rightarrow \mathrm{G}(\zeta)+\hat{\mathrm{i}}_{\mathrm{j}}^{\mathrm{k}}\left(\phi_{\mathrm{k}}-\hat{\phi}_{\mathrm{k}}\right) \tag{11}
\end{equation*}
$$

For the case of an eight node approximation (see Fig. 2),

$$
\begin{align*}
& {\underset{2}{2} 1}_{\Delta}^{\underbrace{}_{1}}\left(\phi_{4}-\hat{\phi}_{4}\right)+{\underset{1}{1} 3}_{\Delta_{7}}\left(\phi_{5}-\hat{\phi}_{5}\right)+ \\
& \Delta_{32}\left(\phi_{6}-\hat{\phi}_{6}\right)+\Delta_{2}\left(\phi_{7}-\hat{\phi}_{7}\right)+\Delta_{4}\left(\phi_{8}-\hat{\phi}_{8}\right) \tag{12}
\end{align*}
$$

In the above equation, $\Delta_{1}$ refers to the intial case of having a constant-valued $G(\zeta)$ defined on $\Gamma$, where $G(\zeta)=$ $\phi_{1}$ for all $z \in \Gamma$, due to having only a single node (\#1) defined on $\Gamma$. Given a specified sequence of nodal point insertion on $\Gamma$, the index notation of $\mathrm{i}, \mathrm{j}, \mathrm{k}$ can be simplified to simply using i , as it is known that node k is to follow node i (in the counterclockwise direction) on $\Gamma$, and k is known by being the kth index term. In the following, it will be assumed that a node installation sequence, $S$, is defined so that node numbers $i, j$ are understood when given node number $k$. Consequently,

$$
\begin{equation*}
G(\zeta)=\sum_{k=1}^{m} \Delta_{s_{k} s_{k+1}}^{k}\left(\phi_{k}-\hat{\phi}_{k}\right) \tag{13}
\end{equation*}
$$

where $s_{k}$ is the kth term of $S, s_{m+1}=s_{1}$, and necessarily $\phi_{1}=$ 0 for the initial case of $k=1$. Equation (13) can be rewritten as,

$$
\begin{equation*}
\mathrm{G}(\zeta)=\sum_{\mathrm{k}=1}^{\mathrm{m}} \Delta^{\mathrm{k}}\left(\phi_{\mathrm{k}}-\hat{\phi}_{\mathrm{k}}\right) \tag{14}
\end{equation*}
$$

where a nodal point installation sequence, $S$, is defined, and nodes i and j are known given node k .

From the above, the extension of a complex variable function $\omega(z)$, defined on $\Gamma$ is given, for $m$ nodes on $\Gamma$, by

$$
\begin{equation*}
\mathrm{G}(\zeta)=\sum_{\mathrm{k}=1}^{\mathrm{m}} \Delta^{\mathrm{k}}\left(\omega_{\mathrm{k}}-\widehat{\omega}_{\mathrm{k}}\right) \tag{15}
\end{equation*}
$$

where $\omega\left(z_{k}\right)=\omega_{k}, k=1,2, \ldots m ;$ and as before, necessarily $\widehat{\omega}_{1}$ $=0$.

## EXTENSION TO THE CVBEM

The CVBEM approximation function can be written as, for $m$ nodes on $\Gamma$,

$$
\begin{equation*}
\widehat{\omega}(z)=\frac{1}{2 \pi i} \int_{\Gamma} \frac{G(\zeta) d \zeta}{\zeta-z}=\frac{1}{2 \pi i} \int_{\Gamma} \frac{\sum_{k=1}^{m} \Delta^{k}\left(\omega_{k}-\hat{\omega}_{k}\right) d \zeta}{\zeta-z}, z \in \Omega \tag{16}
\end{equation*}
$$

or

$$
\begin{equation*}
\widehat{\omega}(\mathrm{z})=\frac{1}{2 \pi \mathrm{i}} \sum_{\mathrm{k}=1}^{\mathrm{m}}\left(\omega_{\mathrm{k}}-\hat{\omega}_{\mathrm{k}}\right) \int_{\Gamma}^{\mathrm{k}} \frac{\Delta \mathrm{~d} \zeta}{\zeta-z}, z \in \Omega \tag{17}
\end{equation*}
$$

For the case of $\omega(z)$ being analytic on $\Gamma$, then $\omega(z)$ is continuous on $\Gamma$ where $G(\zeta) \rightarrow \omega(\zeta)$ on $m \rightarrow \infty$ (and the arclength between successive nodes $\rightarrow 0$ ), and from Schauder's theorem (see Cheney, 1966),

$$
\begin{equation*}
\omega(z)=\frac{1}{2 \pi i} \sum_{\mathrm{k}=1}^{\infty}\left(\omega_{\mathrm{k}}-\hat{\omega}_{\mathrm{k}}\right) \int_{\Gamma}^{\mathrm{k}} \frac{\Delta \mathrm{~d} \zeta}{\zeta-\mathrm{z}}, \mathrm{z} \in \Omega \tag{18}
\end{equation*}
$$

In the above,

## CONCLUSIONS

In this paper, the development of triangular fractals that geometrically sum into an area whose boundary is a function, of a specific type, is used to expand the Complex Variable Boundary Element Method (or CVBEM) into a series. Further research is needed in evaluating convergence and uniqueness properties.

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578
Boundary Elements


Figure 1. $\quad G(\zeta)=\phi_{1}+\Delta_{1}^{2} \phi_{2}+\Delta_{1}^{3} \phi_{3}+\Delta_{2}^{4} \phi_{4}$.


Figure 2. $\quad G(\zeta)=\sum_{k=1}^{8} \Delta_{i}^{k} \phi_{k}$.

